

Advanced Mathematical Methods for Engineers

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$$1) \begin{cases} y'(x) = y(x)^{\frac{3}{2}} \sqrt{y(x)-1} \\ y(0) = K \in \mathbb{R} \end{cases}$$

a) $f(x,y) = f(y) = y^{\frac{3}{2}} \sqrt{y-1} \in C^0(\mathbb{R})$ but it is not differentiable for $y=1$.

For the Peano theorem, we have existence of local solutions $\forall K \in \mathbb{R}$.

Uniqueness can be inferred only for $K \neq 1$ by the Local Cauchy-Lipschitz theorem.

The function f is not globally Lipschitz and so we cannot conclude anything about global solutions.

b) $y=0$ and $y=1$ are solutions in case $K=0$ and $K=1$ respectively.

$y=0$ cannot be intersected by other solutions due to uniqueness of solutions.

Since $y' > 0$ for $y < 0$ and $y > 1$ the solutions are increasing in these sets.

Viceversa they are decreasing on the complementary sets.

The equation is a separable-variable type equation and so we have

$$g_K(y) : \int_K^y \frac{ds}{s^{\frac{3}{2}} \sqrt{s-1}} = t \quad \begin{array}{l} \text{for } K \neq 0 \\ \text{for } K \neq 1 \end{array} \quad ①$$

Hence, we get

b1) for $K < 0$, $\lim_{y \rightarrow 0} g_K(y) = +\infty$

and $\lim_{y \rightarrow -\infty} g_K(y) = T_1 > -\infty$

Hence $\text{dom } y = (T_1, +\infty)$ with

$\lim_{t \rightarrow T_1^+} y(t) = -\infty$ and

$\lim_{t \rightarrow +\infty} y(t) = 0$.

b2) for $K \in (0, 1)$, $\lim_{y \rightarrow 0} g_K(y) = +\infty$ and

$\lim_{y \rightarrow 1} g_K(y) = T_2 > -\infty \Rightarrow$

$\text{dom } y = (T_2, +\infty)$ with $\lim_{t \rightarrow +\infty} y(t) = 0$

and $\lim_{t \rightarrow T_2^+} y(t) = 1$ and

$\lim_{t \rightarrow T_2^+} y'(t) = 0$

b3) for $K > 1$, $\lim_{y \rightarrow +\infty} g_K(y) = T_3 < +\infty$

and $\lim_{y \rightarrow 1} g_K(y) = T_4 > -\infty$

$\Rightarrow \text{dom } y = (T_4, T_3)$, moreover

$\lim_{t \rightarrow T_3^-} y(t) = +\infty$, $\lim_{t \rightarrow T_4^+} y(t) = 1$

and $\lim_{t \rightarrow T_4^+} y'(t) = 0$.

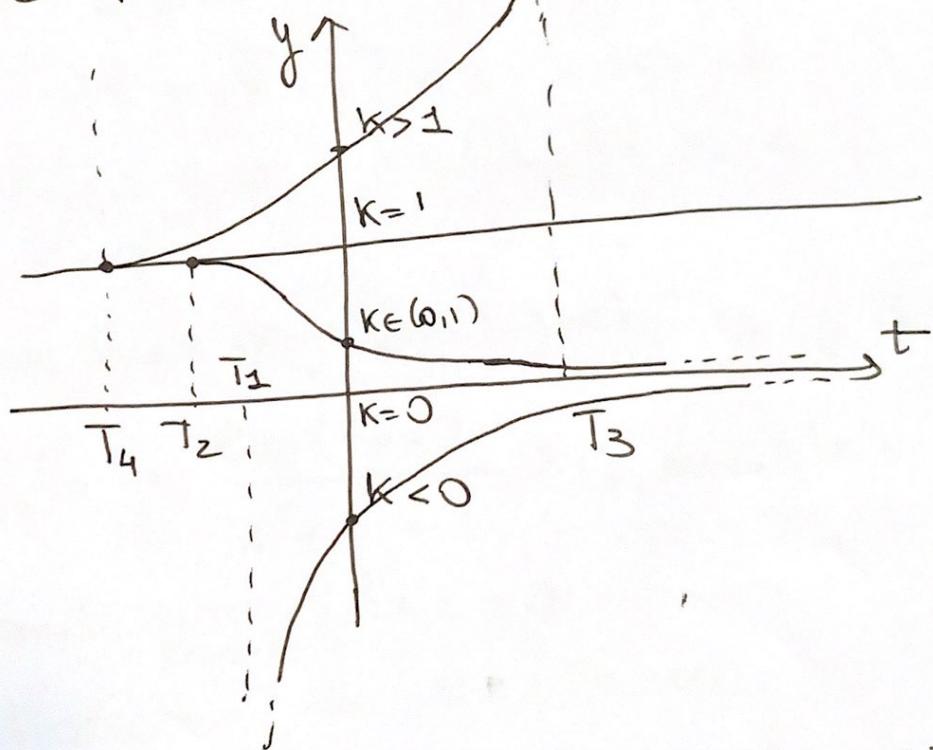
(2)

Moreover, we have

$$y'' = \frac{4y-3}{3\sqrt{(y-1)^2}} \cdot y'$$

and so the flex are on the line $y = \frac{3}{4}$.

The graphs are then the following:



2) Defining $y := x$, we get the ODE system:

$$\begin{cases} x' = y \\ y' = -k(x-2)y - \tan x \end{cases}$$

The matrix of coefficients is:

$$A = \begin{pmatrix} 0 & 1 \\ -1 & 2k \end{pmatrix}$$

(3)

for the linearized system in the origin.

$$\det(A - \lambda \text{Id}) = -\lambda \cdot (2k-\lambda) + 1 = \\ = \lambda^2 - 2k\lambda + 1 = 0$$

$$\Leftrightarrow \lambda_{1,2} = \frac{k \pm \sqrt{k^2 - 1}}{1}$$

$$\text{If } k^2 - 1 \geq 0 \quad \operatorname{Re}(\lambda_{1,2}) = k \pm \sqrt{k^2 - 1} < 0$$

$$\Leftrightarrow k + \sqrt{k^2 - 1} < 0 \Leftrightarrow k < 0$$

$$\text{If } k^2 - 1 < 0 \quad \operatorname{Re}(\lambda_{1,2}) = k < 0$$

$$\Leftrightarrow k < 0$$

\Rightarrow we have asymptotic stability if $k < 0$.

3) $f_m(x) = \frac{\log(1+x)}{x^2 + m^2 + 1} \chi_{[0, m+\sqrt{m}]}(x)$

and $\lim_{m \rightarrow \infty} f_m(x) = 0$.

Moreover $\forall m, \forall x \in [0, +\infty)$

$$|f_m(x)| \leq \frac{\log(1+x)}{x^2 + 1} \in L^1(0, +\infty)$$

\Rightarrow we can apply the Lebesgue dominated convergence theorem and so we get

$$\lim_{m \rightarrow \infty} \int_0^{+\infty} f_m(x) dx = \int_0^{+\infty} \lim_{m \rightarrow \infty} f_m(x) dx = 0.$$

4) We use the separation of variable method:

$$u(x,t) = \varphi(x)\omega(t) \text{ and we get}$$

$$\omega'(t) = \lambda \omega(t) \text{ and}$$

$$\varphi''(x) = \lambda \varphi(x)$$

We first solve the second using the boundary conditions:

$$\varphi(0) = 0, \varphi'(\pi) = 0$$

If $\lambda = \mu^2 \geq 0 \Rightarrow$ we get $\varphi = 0$ and we have only trivial solution

If $\lambda = -\mu^2 < 0 \Rightarrow \lambda_k = -\left(\frac{(2k+1)}{2}\right)^2, k=0,1,\dots$ with corresponding eigenfunctions

$$\varphi_k(x) = \sin\left(\frac{(2k+1)}{2}x\right)$$

\Rightarrow we get

$$u_k(x,t) = c_k \sin\left(\frac{(2k+1)}{2}x\right) e^{-\left(\frac{(2k+1)}{2}\right)^2 t}$$

$$\Rightarrow u(x,t) = \sum c_k \sin\left(\frac{(2k+1)}{2}x\right) e^{-\left(\frac{(2k+1)}{2}\right)^2 t}$$

Finally we get

$$\text{where } c_k = \frac{2}{\pi} \int_0^\pi g(x) \sin\left(\frac{(2k+1)}{2}x\right) dx$$

are the Fourier coefficients of g .

If $g(x) = 2 \sin\left(\frac{3}{2}x\right)$ and we take

$c_1 = 2 \Rightarrow u_1$ satisfies also the initial condition \Rightarrow

$$u(x,t) = 2 e^{-\frac{9}{4}t} \sin\left(\frac{3}{2}x\right) \text{ is the solution.}$$

(5)