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Materials with Memory: free energies & Solution Exponential Decay

- General Framework of the Problem
- Constitutive and Functional Assumptions
- Evolution Problems
- Semigroups Theory: Exponential Decay
- Bäcklund Transformations
- Traveling Wave Solutions

General Framework of the Problem

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key references

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...and many more !!

Quantities of Interest

$$\theta = \theta(\mathbf{x}, t)$$

temperature

$$\theta_0$$

fixed reference temperature

$$u := \theta(\mathbf{x}, t) - \theta_0$$

relative temperature

$$\mathbf{g} := \nabla u(\mathbf{x}, t)$$

temperature gradient

$$\mathbf{q} = \mathbf{q}(\mathbf{x}, t)$$

heat flux

$$e = e(\mathbf{x}, t)$$

internal energy

$$\alpha_0 = \text{const}$$

energy relaxation coefficient

$$k = k(\mathbf{x}, t)$$

heat flux relaxation function

$$k_0(\mathbf{x}) := k(\mathbf{x}, 0)$$

initial heat flux relaxation function

furthermore constitutive assumptions

Constitutive Assumptions

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$$e(\mathbf{x}, t) = \alpha_0 u(\mathbf{x}, t) \quad , \quad \alpha_0 \text{ const}$$

$$\mathbf{q}(\mathbf{x}, t) = - \int_0^\infty k(\mathbf{x}, \tau) \nabla u(\mathbf{x}, t - \tau) d\tau$$

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the heat flux can be equivalently written in the form:

$$\mathbf{q}(\mathbf{x}, t) = \int_0^\infty \dot{k}(\mathbf{x}, \tau) \bar{\mathbf{g}}^t(\mathbf{x}, \tau) d\tau . \quad (-1)$$

where $\mathbf{g} := \nabla u$, and $\bar{\mathbf{g}}^t(\mathbf{x}, \tau) := \int_{t-\tau}^t \nabla u(\mathbf{x}, s) ds$

$\bar{\mathbf{g}}^t(\mathbf{x}, \tau)$... **integrated history of the temperature gradient**

Functional Assumptions

- $\dot{k} \in L^1(\mathbb{R}^+) \cap L^2(\mathbb{R}^+) \quad \forall \mathbf{x} \in \Omega$
- $k \in L^1(\mathbb{R}^+) \quad \forall \mathbf{x} \in \Omega$
- $k(\mathbf{x}, t) = k_0(\mathbf{x}) + \int_0^t \dot{k}(\mathbf{x}, s) \, ds$
- $k(\mathbf{x}, \infty) := \lim_{t \rightarrow \infty} k(\mathbf{x}, t) = 0$

fading memory property $\forall \mathbf{x} \in \Omega$

$$\forall \varepsilon > 0 \exists \tilde{a} = a(\varepsilon, \bar{\mathbf{g}}^t) \in \mathbb{R}^+ \text{ s.t. } \forall a > \tilde{a} \Rightarrow \left| \int_0^\infty \dot{k}(s+a) \bar{\mathbf{g}}^t(s) \, ds \right| < \varepsilon$$

...i.e., no heat flux when, at infinity, equilibrium is reached!

The evolution problem

$$\left\{ \begin{array}{ll} \alpha_0 u_t = -\nabla \cdot \mathbf{q} + r(\mathbf{x}, t), & r(\mathbf{x}, t) \text{ heat supply} \\ \mathbf{q}(\mathbf{x}, t) = \int_0^\infty \dot{k}(\mathbf{x}, \tau) \overline{\nabla u}^t(\mathbf{x}, \tau) d\tau & \text{in } Q_T = \Omega \times [0, T) \end{array} \right.$$

initial and boundary conditions

Problem 1

no heat flux through the body boundary (Neumann cond.s)

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no heat flux through the body boundary (Neumann cond.s)

$$\begin{aligned} u(\mathbf{x}, 0) &= u_0(\mathbf{x}) & \forall \mathbf{x} \in \Omega \\ \mathbf{q}(\mathbf{x}, t) \cdot \mathbf{n}(\mathbf{x})|_{\partial\Omega} &= 0, & \forall t \in [0, T) \end{aligned}$$

where $\mathbf{n}(\mathbf{x})$ denotes the outward unit normal to the (smooth) boundary $\partial\Omega$ at $\mathbf{x} \in \partial\Omega$ and, in addition, say $r(\mathbf{x}, t) = 0$, i.e. no heat supply.

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initial and boundary conditions

Problem 2

fixed temperature at the body boundary (Dirichlet cond.s)

$$\begin{aligned} u(\mathbf{x}, 0) &= u_0(\mathbf{x}) & \forall \mathbf{x} \in \Omega \\ u(\mathbf{x}, t)|_{\partial\Omega} &= 0, & \forall t \in [0, T) \end{aligned}$$

at the boundary of Ω , namely $\forall \mathbf{x} \in \partial\Omega$, and, in addition, say $r(\mathbf{x}, t) = 0$, i.e. no heat supply.

Semigroups Theory

Semigroups Theory

- ▷ write **evolution problem** as a dynamical system

$$\begin{cases} \dot{u}(x, t) = \nabla \cdot \int_0^{+\infty} \dot{k}(\tau) \bar{\mathbf{g}}^t(\tau) d\tau \\ \dot{\bar{\mathbf{g}}}^t(\tau) = \nabla u(t) - \nabla u(t - \tau) \end{cases}$$

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- ▷ **energy inner product** ...

$$\langle \chi_1(t), \chi_2(t) \rangle = \frac{1}{2} \int_{\Omega} \alpha(\mathbf{x}) u_1(\mathbf{x}, t) u_2(\mathbf{x}, t) d\mathbf{x} - \frac{1}{2} \int_{\Omega} \int_0^{\infty} \dot{k}(\tau) \bar{\mathbf{g}}_1^t(\tau) \cdot \bar{\mathbf{g}}_2^t(\tau) d\tau d\mathbf{x}$$

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- ▷ ... and induced norm

$$\|\chi(t)\|^2 = \frac{1}{2} \int_{\Omega} \alpha(\mathbf{x}) u^2(\mathbf{x}, t) d\mathbf{x} - \frac{1}{2} \int_{\Omega} \int_0^{\infty} \dot{k}(\tau) [\bar{\mathbf{g}}^t(\tau)]^2 d\tau d\mathbf{x}$$

Exponential Decay

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▷ recall Graffi's free energy

$$\zeta_G(u, \bar{\mathbf{g}}^t) := \frac{\alpha_0}{2} [u(\mathbf{x}, t)]^2 - \frac{1}{2} \int_0^\infty k'(\tau) [\bar{\mathbf{g}}^t(\mathbf{x}, \tau)]^2 d\tau$$

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- ▷ Graffi's free energy satisfies

$$\frac{\partial \zeta_G}{\partial t}(\mathbf{x}, t) - \frac{\partial e}{\partial t}(\mathbf{x}, t) u(\mathbf{x}, t) + \mathbf{q}(\mathbf{x}, t) \cdot \mathbf{g}(\mathbf{x}, t) = -\frac{1}{2} \int_0^\infty k''(\tau) [\bar{\mathbf{g}}^t(\mathbf{x}, \tau)]^2 d\tau$$

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- ▷ introduce the energy

$$\mathcal{E}(t) = \frac{1}{2} \int_\Omega \alpha_0 [u(\mathbf{x}, t)]^2 d\mathbf{x} - \frac{1}{2} \int_\Omega \int_0^\infty k'(\tau) [\bar{\mathbf{g}}^t(\mathbf{x}, \tau)]^2 d\tau d\mathbf{x}$$

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▷ under the assumption

$$0 < -\dot{k}(\tau) \leq \delta \ddot{k}(\tau) \quad , \quad \dot{k}(0) \text{ bounded}$$

since

$$\int_{\Omega} \dot{\zeta}_G(u, \bar{\mathbf{g}}^t) d\mathbf{x} = -\frac{1}{2} \int_{\Omega} \int_0^{\infty} \ddot{k}(\mathbf{x}, \tau) [\bar{\mathbf{g}}^t(\tau)]^2 d\tau d\mathbf{x} \leq 0$$

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▷ exponential decay follows

$$\frac{d}{dt} \mathcal{E}(t) \leq 0 \implies 0 \leq \mathcal{E}(t) \leq \mathcal{E}(0)$$

Bäcklund Transformations

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► given two Evolution Equations, K and G are C^∞ -vector field

$$\begin{aligned} u_t &= K(u), & K : M_1 &\rightarrow TM_1, & u : (x, t) \in \mathbb{R}^n \times \mathbb{R} &\rightarrow u(x, t) \in \mathbb{R}^r \\ v_t &= G(v), & G : M_2 &\rightarrow TM_2, & v : (x, t) \in \mathbb{R}^n \times \mathbb{R} &\rightarrow v(x, t) \in \mathbb{R}^n \end{aligned}$$

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- **Bäcklund Transformations** [Fuchssteiner, Nonlin. An. TMA, 3, 849,80]. let,
in turn, $u(x, t)$ and $v(x, t)$ be solutions to the two evolution
equations

$$\text{if } \boxed{B(u, v)|_{t=0} = 0 \Rightarrow B(u, v) = 0 \quad \forall t}$$

then $\boxed{B(u, v) = 0}$ is a Bäcklund Transformations

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- ▶ if $u_t = K(u)$ admits a Recursion Operator Φ it can be written

$$u_t = \Phi(u)[u_x]$$

Bäcklund Chart

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- ▶ Bäcklund Chart: net of links connecting evolution equations

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- ▶ Example 1 [Cole(1950),Hopf(1950)].

$$\boxed{u_t = u_{xx} + 2uu_x} \xleftrightarrow{BT} \boxed{v_t = v_{xx}}$$

$$B(u, v) = 0 \text{ reads } \boxed{v_x - uv = 0}$$

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- ▶ **Example 1** [Cole(1950),Hopf(1950)].

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- ▶ **Example 2**

$$\boxed{u_t = u_{xxx} + 6uu_x} \xleftrightarrow{BT} \boxed{v_t = v_{xxx} - 6v^2v_x}$$

$$B(u, v) = 0 \text{ reads } \boxed{u + v_x + v^2 = 0}$$

...Latest Bäcklund Chart

[S. C., 2008]

...Latest Bäcklund Chart

[S. C., 2008]

Heat Conduction with memory:

...Latest Bäcklund Chart

[S. C., 2008]

Heat Conduction with memory:

- Consider a 1-dimesional heat conductor with memory

$$\alpha_0 v_t = - \left[\int_0^t \dot{k}(\tau) \overline{\nabla v}^t(\mathbf{x}, \tau) d\tau \right]_x$$

which ... after a Cole-Hopf transformations is related to

$$u_{tt} = k_0 [u_{xx} + 2uu_x] \quad , \quad k_0 = k(0)$$

... look for solutions of the latter ...

Traveling Wave Solutions

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► Inspection shows [S. C., Vietri 2008]

$$u(x, t) = \frac{k_0 - c^2}{k_0(x \pm ct + C)}, \quad \forall C \in \mathbb{R}, \quad \forall c \in \mathbb{R}^+.$$

$$c^2 = k_0 \implies \text{trivial zero solution.}$$

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$c^2 = k_0 \implies$ trivial zero solution.

- further solutions are $\forall c \in \mathbb{R}^+ \setminus \{\sqrt{k_0}\}, \forall c_1, c_2 \in \mathbb{R},$

$$u(x \pm ct) = - \frac{c_1 e^{\frac{k_0}{(c^2 - k_0)}(x \pm ct)} - c_2 e^{-\frac{k_0}{(c^2 - k_0)}(x \pm ct)}}{c_1 e^{\frac{k_0}{(c^2 - k_0)}(x \pm ct)} + c_2 e^{-\frac{k_0}{(c^2 - k_0)}(x \pm ct)}},$$

which also represent *traveling wave* solutions of the nonlinear wave equation

$$u_{tt} = k_0 [u_{xx} + 2uu_x] \quad , \quad k_0 = k(0)$$

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