

Vector calculus MA3VC 2014–15 — Assignment 1

SOLUTIONS

(Exercise 1) Compute a scalar potential φ for the vector field $\vec{\mathbf{F}} = yz(z\hat{\mathbf{j}} + y\hat{\mathbf{k}}) + \cos 2\pi x\hat{\mathbf{i}}$.
Is $\vec{\mathbf{F}}$ solenoidal, irrotational? Does $\vec{\mathbf{F}}$ allow a vector potential?

It is easy to find that the scalar potentials of $\vec{\mathbf{F}}$ are the scalar fields $\varphi = \frac{1}{2\pi} \sin 2\pi x + \frac{1}{2}y^2z^2 + \lambda$, where λ is a real constant:

$$\begin{aligned} \frac{\partial \varphi}{\partial x} = \cos 2\pi x &\Rightarrow \varphi(x, y, z) = \frac{1}{2\pi} \sin 2\pi x + f(y, z) && \text{for some two-dimensional scalar field } f, \\ \frac{\partial \varphi}{\partial y} = yz^2 &\Rightarrow \frac{\partial(\frac{1}{2\pi} \sin 2\pi x + f(y, z))}{\partial y} = \frac{\partial f(y, z)}{\partial y} = yz^2 &\Rightarrow f = \frac{1}{2}y^2z^2 + g(z) \\ &\Rightarrow \varphi = \frac{1}{2\pi} \sin 2\pi x + \frac{1}{2}y^2z^2 + g(z) && \text{for some real function } g, \\ \frac{\partial \varphi}{\partial z} = y^2z &\Rightarrow \frac{\partial(\frac{1}{2\pi} \sin 2\pi x + \frac{1}{2}y^2z^2 + g(z))}{\partial z} = y^2z + \frac{\partial g(z)}{\partial z} = y^2z &\Rightarrow \frac{\partial g(z)}{\partial z} = 0 \\ &\Rightarrow \varphi = \frac{1}{2\pi} \sin 2\pi x + \frac{1}{2}y^2z^2 + \lambda. \end{aligned}$$

To verify that the scalar potential is correct, it is sufficient to check that $\vec{\nabla}\varphi = \vec{\mathbf{F}}$.

Since $\vec{\mathbf{F}}$ is conservative, by the box in Section 1.5 or by the identity $\vec{\nabla} \times (\vec{\nabla}\varphi) = \vec{\mathbf{0}}$, $\vec{\mathbf{F}}$ is irrotational.

The divergence of $\vec{\mathbf{F}}$ is not zero: $\vec{\nabla} \cdot \vec{\mathbf{F}} = -2\pi \sin 2\pi x + z^2 + y^2 \neq 0$, so $\vec{\mathbf{F}}$ is not solenoidal. This implies that $\vec{\mathbf{F}}$ does not admit a vector potential, again by the box in Section 1.5 or by the identity $\vec{\nabla} \cdot (\vec{\nabla} \times \vec{\mathbf{A}}) = 0$.

(Exercise 2) Let $\vec{\mathbf{F}}$ be a vector field with scalar potential φ , and let $\vec{\mathbf{G}}$ be a vector field with scalar potential ψ . Prove the following identity: $2\vec{\mathbf{F}} \cdot \vec{\mathbf{G}} = \Delta(\varphi\psi) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}}$.

The identity to be proved is nothing else than the product rule (32) for the Laplacian in disguise. We can either use the known vector identities (simpler, *i*) or expand in partial derivatives (more complicated, *ii*).¹

(Version *i*) We use three tools:

- The definition of scalar potential, namely $\vec{\mathbf{F}} = \vec{\nabla}\varphi$ and $\vec{\mathbf{G}} = \vec{\nabla}\psi$;
- Identity (22) in the notes, which gives $\Delta\varphi = \vec{\nabla} \cdot (\vec{\nabla}\varphi) = \vec{\nabla} \cdot \vec{\mathbf{F}}$ and $\Delta\psi = \vec{\nabla} \cdot (\vec{\nabla}\psi) = \vec{\nabla} \cdot \vec{\mathbf{G}}$;
- The product rule (32) for the Laplacian.

These identities together lead to

$$\Delta(\varphi\psi) \stackrel{(32)}{=} (\Delta\varphi)\psi + 2\vec{\nabla}\varphi \cdot \vec{\nabla}\psi + (\Delta\psi)\varphi \stackrel{(22)}{=} \vec{\nabla} \cdot (\vec{\nabla}\varphi)\psi + 2\vec{\nabla}\varphi \cdot \vec{\nabla}\psi + \vec{\nabla} \cdot (\vec{\nabla}\psi)\varphi = (\vec{\nabla} \cdot \vec{\mathbf{F}})\psi + 2\vec{\mathbf{F}} \cdot \vec{\mathbf{G}} + (\vec{\nabla} \cdot \vec{\mathbf{G}})\varphi.$$

Rearranging for $2\vec{\mathbf{F}} \cdot \vec{\mathbf{G}}$ immediately gives the desired result.

(Version *ii*) If we use the expansion in components, we need to use twice the product rule for partial derivatives (8), together with the definitions of Laplacian (20), divergence (17) and scalar potentials $\vec{\mathbf{F}} = \vec{\nabla}\varphi$, $\vec{\mathbf{G}} = \vec{\nabla}\psi$. The right-hand side of the identity can be expanded as follows:

$$\begin{aligned} \Delta(\varphi\psi) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}} &\stackrel{(17),(20)}{=} \frac{\partial^2(\varphi\psi)}{\partial x^2} + \frac{\partial^2(\varphi\psi)}{\partial y^2} + \frac{\partial^2(\varphi\psi)}{\partial z^2} - \psi \left(\frac{\partial F_1}{\partial x} + \frac{\partial F_2}{\partial y} + \frac{\partial F_3}{\partial z} \right) - \varphi \left(\frac{\partial G_1}{\partial x} + \frac{\partial G_2}{\partial y} + \frac{\partial G_3}{\partial z} \right) \\ &\stackrel{(8)}{=} \frac{\partial}{\partial x} \left(\frac{\partial \varphi}{\partial x} \psi + \varphi \frac{\partial \psi}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\partial \varphi}{\partial y} \psi + \varphi \frac{\partial \psi}{\partial y} \right) + \frac{\partial}{\partial z} \left(\frac{\partial \varphi}{\partial z} \psi + \varphi \frac{\partial \psi}{\partial z} \right) \\ &\quad - \psi \left(\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} + \frac{\partial^2 \varphi}{\partial z^2} \right) - \varphi \left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) \\ &\stackrel{(8)}{=} \frac{\partial^2 \varphi}{\partial x^2} \psi + 2 \frac{\partial \varphi}{\partial x} \frac{\partial \psi}{\partial x} + \varphi \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \psi + 2 \frac{\partial \varphi}{\partial y} \frac{\partial \psi}{\partial y} + \varphi \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \varphi}{\partial z^2} \psi + 2 \frac{\partial \varphi}{\partial z} \frac{\partial \psi}{\partial z} + \varphi \frac{\partial^2 \psi}{\partial z^2} \end{aligned}$$

¹ A nice alternative solution I found in some of the assignments is the following (similar to *i* but slightly more complicated):

$$\begin{aligned} 2\vec{\mathbf{F}} \cdot \vec{\mathbf{G}} &= \vec{\mathbf{F}} \cdot \vec{\mathbf{G}} + \vec{\mathbf{F}} \cdot \vec{\mathbf{G}} = \vec{\nabla}\varphi \cdot \vec{\nabla}\psi + \vec{\mathbf{F}} \cdot \vec{\nabla}\psi \stackrel{(28)}{=} \vec{\nabla} \cdot (\varphi\vec{\mathbf{G}}) - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}} + \vec{\nabla} \cdot (\psi\vec{\mathbf{F}}) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} = \vec{\nabla} \cdot (\varphi\vec{\mathbf{G}} + \psi\vec{\mathbf{F}}) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}} \\ &= \vec{\nabla} \cdot (\varphi\vec{\nabla}\psi + \psi\vec{\nabla}\varphi) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}} \stackrel{(26)}{=} \vec{\nabla} \cdot (\vec{\nabla}(\varphi\psi)) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}} \stackrel{(22)}{=} \Delta(\varphi\psi) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}}. \end{aligned}$$

$$\begin{aligned}
& -\psi\left(\frac{\partial^2\varphi}{\partial x^2} + \frac{\partial^2\varphi}{\partial y^2} + \frac{\partial^2\varphi}{\partial z^2}\right) - \varphi\left(\frac{\partial^2\psi}{\partial x^2} + \frac{\partial^2\psi}{\partial y^2} + \frac{\partial^2\psi}{\partial z^2}\right) \\
&= 2\frac{\partial\varphi}{\partial x}\frac{\partial\psi}{\partial x} + 2\frac{\partial\varphi}{\partial y}\frac{\partial\psi}{\partial y} + 2\frac{\partial\varphi}{\partial z}\frac{\partial\psi}{\partial z} \\
&= 2\vec{\nabla}\varphi \cdot \vec{\nabla}\psi \\
&= 2\vec{\mathbf{F}} \cdot \vec{\mathbf{G}}.
\end{aligned}$$

(Exercise 3) Demonstrate the identity in Ex. (2) for the vector field $\vec{\mathbf{F}}$ in Ex. (1) and the scalar field $\psi = y^3$.

We compute all the terms appearing in the identity (λ can be fixed to 0):

$$\begin{aligned}
\varphi &= \frac{1}{2\pi} \sin 2\pi x + \frac{1}{2}y^2z^2 + \lambda && \text{from Exercise (1),} \\
\psi &= y^3, \\
\vec{\mathbf{F}} &= \cos 2\pi x \hat{\mathbf{i}} + yz(z\hat{\mathbf{j}} + y\hat{\mathbf{k}}), \\
\vec{\mathbf{G}} &= \vec{\nabla}\psi = 3y^2\hat{\mathbf{j}}, \\
\vec{\nabla} \cdot \vec{\mathbf{F}} &= \frac{\partial(\cos 2\pi x)}{\partial x} + \frac{\partial(yz^2)}{\partial y} + \frac{\partial(y^2z)}{\partial z} = -2\pi \sin 2\pi x + z^2 + y^2, \\
\vec{\nabla} \cdot \vec{\mathbf{G}} &= \frac{\partial(3y^2)}{\partial y} = 6y, \\
\psi(\vec{\nabla} \cdot \vec{\mathbf{F}}) &= \psi\Delta\varphi = -2\pi y^3 \sin 2\pi x + y^3z^2 + y^5, \\
\varphi(\vec{\nabla} \cdot \vec{\mathbf{G}}) &= \varphi\Delta\psi = \frac{3}{\pi}y \sin 2\pi x + 3y^3z^2 + 6y\lambda, \\
\vec{\mathbf{F}} \cdot \vec{\mathbf{G}} &= 3y^3z^2, \\
\varphi\psi &= \frac{1}{2\pi}y^3 \sin 2\pi x + \frac{1}{2}y^5z^2 + y^3\lambda \\
\Delta(\varphi\psi) &= \frac{\partial^2(\frac{1}{2\pi}y^3 \sin 2\pi x)}{\partial x^2} + \frac{\partial^2(\frac{1}{2\pi}y^3 \sin 2\pi x + \frac{1}{2}y^5z^2 + y^3\lambda)}{\partial y^2} + \frac{\partial^2(\frac{1}{2}y^5z^2)}{\partial z^2} \\
&= -2\pi y^3 \sin 2\pi x + \frac{3}{\pi}y \sin 2\pi x + 10y^3z^2 + 6y\lambda + y^5, \\
LHS &= 2\vec{\mathbf{F}} \cdot \vec{\mathbf{G}} = 6y^3z^2, \\
RHS &= \Delta(\varphi\psi) - \psi\vec{\nabla} \cdot \vec{\mathbf{F}} - \varphi\vec{\nabla} \cdot \vec{\mathbf{G}} \\
&= -2\pi y^3 \sin 2\pi x + \frac{3}{\pi}y \sin 2\pi x + 10y^3z^2 + 6y\lambda + y^5 \\
&\quad - \left(-2\pi y^3 \sin 2\pi x + y^3z^2 + y^5\right) - \left(\frac{3}{\pi}y \sin 2\pi x + 3y^3z^2 + 6y\lambda\right) = 6y^3z^2.
\end{aligned}$$

The left-hand side (LHS) and the right-hand side (RHS) of the identity coincide, so the identity is demonstrated.

(Exercise 4) Consider the cylinder C with radius 2, axis of rotation on the x -axis and flat faces lying in the planes $\{x = 0\}$ and $\{x = 10\}$. Write C in the form $C = \{\vec{\mathbf{r}} \in \mathbb{R}^3 \text{ s.t. } \dots\}$ and compute the outward pointing unit normal vector field $\hat{\mathbf{n}}$ defined on the boundary of C .

The cylinder can be written either as an open set

$$C_o = \{\vec{\mathbf{r}} \in \mathbb{R}^3 \text{ such that } 0 < x < 10, y^2 + z^2 < 4\},$$

or as a closed set

$$C_c = \{\vec{\mathbf{r}} \in \mathbb{R}^3 \text{ such that } 0 \leq x \leq 10, y^2 + z^2 \leq 4\}.$$

The two flat faces lie in the planes $\{x = 0\}$ and $\{x = 10\}$, so their normals must be either $\hat{\mathbf{n}} = \hat{\mathbf{i}}$ or $\hat{\mathbf{n}} = -\hat{\mathbf{i}}$. (If this is not clear from geometric intuition, we note that these planes are level sets for the scalar field $f(\vec{\mathbf{r}}) = x$, whose gradient is $\vec{\nabla}f = \hat{\mathbf{i}}$ and has already unit length.) With simple geometric consideration, since $\hat{\mathbf{n}}$ must point outward, it is clear that $\hat{\mathbf{n}} = -\hat{\mathbf{i}}$ on $\{x = 0\}$ and $\hat{\mathbf{n}} = \hat{\mathbf{i}}$ on $\{x = 10\}$.

The side has equation $\{y^2 + z^2 = 4\}$, which is a level set of the field $f(\vec{\mathbf{r}}) = y^2 + z^2$. Its gradient is $\vec{\nabla}f = 2y\hat{\mathbf{j}} + 2z\hat{\mathbf{k}}$, which has length $|\vec{\nabla}f| = 2\sqrt{y^2 + z^2}$. Thus, as in the exercises seen in class or in Example 1.33, $\hat{\mathbf{n}} = \pm\vec{\nabla}f/|\vec{\nabla}f| = \pm(y\hat{\mathbf{j}} + z\hat{\mathbf{k}})/\sqrt{y^2 + z^2}$. Since $\hat{\mathbf{n}}$ is “outward pointing”, it must point away from the x -axis, e.g. it must satisfy $\hat{\mathbf{n}}(\hat{\mathbf{i}} + \hat{\mathbf{j}}) = \hat{\mathbf{j}}$, so we choose its sign as

$$\hat{\mathbf{n}}(\vec{\mathbf{r}}) = \frac{y\hat{\mathbf{j}} + z\hat{\mathbf{k}}}{\sqrt{y^2 + z^2}}.$$