## Classical computational methods, Part II

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#### **Some examples**

#### A second order ODE

$$-u''(x) = f(x)$$

Solution can be explicitly determined (closed form solution)

$$u'(x) = u'(x_0) + \int_{x(0)}^x u''(t) \, dt = u'(x_0) - \int_{x(0)}^x f(t) \, dt$$

$$u(x) = u(x_0) + \int_{x(0)}^{x} u'(t) dt$$

In general

$$u(x) = \alpha + \beta x - \int \int f(t) \, dt$$

One dimensional convection equation (PDE)

$$\frac{\partial u}{\partial t}(x,t) - \frac{\partial u}{\partial x}(x,t) = 0$$

Closed form solution

$$u(x,t) = w(x+t)$$



One dimensional wave equation

$$\frac{\partial^2 u}{\partial t^2}(x,t) - c^2 \frac{\partial^2 u}{\partial x^2}(x,t) = 0$$

Closed form solution

$$u(x,t) = w_1(x+ct) + w_2(x-ct)$$

Proof



One dimensional heat equation

$$\frac{\partial u}{\partial t}(x,t) - \frac{\partial^2 u}{\partial x^2}(x,t) = 0, \qquad x \in (0,1), \ t > 0$$

Closed form solution

$$u(x,t) = \sum_{j=1}^{\infty} u_{0,j} e^{-(j\pi)^2 t} \sin(j\pi x),$$

where  $u_0(x) = u(x, 0)$  is the initial datum and

$$u_{0,j} = 2 \int_0^1 u_0(x) \sin(j\pi x) \, dx$$

#### Convection equation

$$\frac{\partial u}{\partial t} + \operatorname{div}(\vec{\beta} u) = 0$$

First order linear equation.

N.B.: divergence operator div  $\vec{v} = \sum_{i=1}^{d} \frac{\partial v_i}{\partial x_i}$ 

This equation states the mass conservation of a body occupying a region  $\Omega \in \mathbb{R}^d$ , with density u and velocity  $\vec{\beta}$ 

Laplace/Beltrami/Poisson equation

 $-\Delta u = f$ 

Second order linear equation.

N.B.: Laplace operator  $\Delta v = \sum_{i=1}^{d} \frac{\partial^2 v}{\partial x_i^2}$ 

This equation states the diffusion of a homogeneous and isotropic fluid occupying a region  $\Omega \in \mathbb{R}^d$ , as well as the vertical displacement of an elastic membrane. Fundamental equation for several models.

#### Heat equation

$$\frac{\partial u}{\partial t} - \Delta u = f$$

Second order linear equation

#### Wave equation

$$\frac{\partial^2 u}{\partial t^2} - \Delta u = 0$$

Second order linear equation

Burgers equation (d = 1)

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0$$

First order quasi-linear equation

Viscous Burgers equation (d = 1)

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \varepsilon \frac{\partial^2 u}{\partial x^2}, \qquad \varepsilon > 0$$

Second order semi-linear equation

### Plan of the Course

Classification of Partial Differential Equations (PDE)

#### Elliptic PDE's

- Finite differences
- Finite elements
- Where the theory is elegant and complete. . .

From elliptic to hyperbolic PDE's

- Convection-diffusion equation
- Finite differences, upwind
- Integrating along the characteristics
- Stabilization of finite elements

#### Parabolic equation

- Heat equation: space semidiscretization and evolution in time
  Stability of θ-method
- Examples with Matlab
- Conclusions and comments
- Questions and answers

## **Classification of (linear) PDE's**

The case of two variables (can be generalized)

$$Lu \equiv \left(A\frac{\partial^2 u}{\partial x_1^2} + B\frac{\partial^2 u}{\partial x_1 \partial x_2} + C\frac{\partial^2 u}{\partial x_2^2}\right) + L.O.T.$$

Matrix associated to quadratic form

$$QF = \left(\begin{array}{cc} A & \frac{1}{2}B\\ \frac{1}{2}B & C \end{array}\right)$$

Note: A, B, and C might be functions themselves.

Compute eigenvalues  $\lambda_i$  of QF

- ► *Elliptic* equation:  $\lambda_1 \lambda_2 > 0$
- ▶ *Parabolic* equation:  $\lambda_1 \lambda_2 = 0$
- Hyperbolic equation:  $\lambda_1 \lambda_2 < 0$

With the notation of quadratic forms: *definite* form, *semidefinite* form, *indefinite* form, respectively.

Consider operator

$$\mathcal{L}u \equiv A \frac{\partial^2 u}{\partial x_1^2} + B \frac{\partial^2 u}{\partial x_1 \partial x_2} + C \frac{\partial^2 u}{\partial x_2^2} = 0$$

and look for change of variables

$$\xi = \alpha x_2 + \beta x_1, \quad \eta = \gamma x_2 + \delta x_1$$

so that  $\mathcal{L}u$  is a multiple of  $\frac{\partial^2 u}{\partial \xi \partial \eta}$  (see wave equation)

$$\mathcal{L}u = (A\beta^2 + B\alpha\beta + C\alpha^2)\frac{\partial^2 u}{\partial\xi^2} + (A\delta^2 + B\gamma\delta + C\gamma^2)\frac{\partial^2 u}{\partial\eta^2} + (2A\beta\delta + B(\alpha\delta + \beta\gamma) + 2C\alpha\gamma)\frac{\partial^2 u}{\partial\xi\partial\eta}$$

$$\mathcal{L}u = (A\beta^{2} + B\alpha\beta + C\alpha^{2})\frac{\partial^{2}u}{\partial\xi^{2}} + (A\delta^{2} + B\gamma\delta + C\gamma^{2})\frac{\partial^{2}u}{\partial\eta^{2}} + (2A\beta\delta + B(\alpha\delta + \beta\gamma) + 2C\alpha\gamma)\frac{\partial^{2}u}{\partial\xi\partial\eta}$$

If A = C = 0, trivial. Suppose  $A \neq 0$ ; we want

$$A\beta^2 + B\alpha\beta + C\alpha^2 = 0, \quad A\delta^2 + B\gamma\delta + C\gamma^2 = 0$$

When  $\alpha \gamma \neq 0$ , divide first equation by  $\alpha^2$ , second one by  $\gamma^2$  and solve for  $\beta/\alpha$  and  $\delta/\gamma$ , resp.

$$\beta/\alpha = (2A)^{-1}(-B \pm \sqrt{\Delta}), \quad \delta/\gamma = (2A)^{-1}(-B \pm \sqrt{\Delta})$$

 $\Delta = B^2 - 4AC$ 

Hyperbolic case

 $\xi = \alpha x_2 + \beta x_1, \quad \eta = \gamma x_2 + \delta x_1$ 

 $\beta/\alpha = (2A)^{-1}(-B \pm \sqrt{\Delta}), \quad \delta/\gamma = (2A)^{-1}(-B \pm \sqrt{\Delta})$ 

For nonsingular change of variables,  $\Delta$  must be positive

$$\alpha = \gamma = 2A, \quad \beta = -B + \sqrt{\Delta}, \delta = -B - \sqrt{\Delta}$$

$$\mathcal{L}u = -4A(B^2 - 4AC)\frac{\partial^2 u}{\partial\xi\partial\eta}$$

As before, solution has the form  $u = p(\xi) + q(\eta)$  and the lines  $\xi = constant$  and  $\eta = constant$  are called *characteristics*.

Actually, when  $x_1 = t$  and  $x_2 = x$ , the change of variables

$$x' = x - \frac{B}{2A}t, \quad t' = t$$

maps our hyperbolic operator  $(A \neq 0)$  to a multiple of wave equation



Hence,  $\mathcal{L}$  is a wave operator in a frameset moving at speed -B/(2A).

Parabolic case

$$\mathcal{L}u = (A\beta^{2} + B\alpha\beta + C\alpha^{2})\frac{\partial^{2}u}{\partial\xi^{2}} + (A\delta^{2} + B\gamma\delta + C\gamma^{2})\frac{\partial^{2}u}{\partial\eta^{2}} + (2A\beta\delta + B(\alpha\delta + \beta\gamma) + 2C\alpha\gamma)\frac{\partial^{2}u}{\partial\xi\partial\eta}$$

For  $\beta/\alpha = -B/(2A)$  coefficient of  $\frac{\partial^2 u}{\partial \xi^2}$  vanishes But B/(2A) = 2C/B, so coefficient of  $\frac{\partial^2 u}{\partial \xi \partial \eta}$  is zero as well

Everything can be written as a multiple of  $\frac{\partial^2 u}{\partial \eta^2}$ 

In conclusion, in the parabolic case, the change of variables

 $\xi = 2Ax_2 - Bx_1, \quad \eta = x_1$ 

maps the equation to

$$A\frac{\partial^2 u}{\partial \eta^2} = 0$$

which has the general solution

 $u = p(\xi) + \eta q(\xi)$ 

One family of characteristics  $\xi = constant$ 

#### Elliptic case

No choice of parameters makes coefficients of  $\frac{\partial^2 u}{\partial \xi^2}$  and  $\frac{\partial^2 u}{\partial \eta^2}$  vanish In this case change of variables

$$\xi = \frac{2Ax_2 - Bx_1}{\sqrt{4AC - B^2}}, \quad \eta = x_1$$

#### maps equation to

$$A\left(\frac{\partial^2 u}{\partial\xi^2} + \frac{\partial^2 u}{\partial\eta^2}\right) = 0$$

No family of characteristics (infinite speed of propagation, no discontinuities allowed)

#### Final examples

- Laplace equation: elliptic
- ► Wave equation: hyperbolic
- Heat equation: parabolic
- Convection-diffusion equation:

 $\frac{\partial u}{\partial t} - \varepsilon \Delta u + \operatorname{div}(\vec{\beta}u) = 0$ 

parabolic, degenerating to hyperbolic as  $\varepsilon$  tends to zero.

### End of Part I

### **Closed form of 1D wave equation solution**

Change of variables

 $y = x + ct, \quad z = x - ct, \qquad u(x,t) = w(y,z)$   $\frac{\partial^2 u}{\partial t^2} = c^2 \left( \frac{\partial^2 w}{\partial y^2} - 2 \frac{\partial^2 w}{\partial y \partial z} + \frac{\partial^2 w}{\partial z^2} \right)$   $\frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2}$   $\frac{\partial^2 w}{\partial y \partial z} = 0 \Rightarrow w = w_1(y) + w_2(z) = w_1(x + ct) + w_2(x - ct)$